

CMB resonant scattering as a cosmological tool

Kaustuv Basu

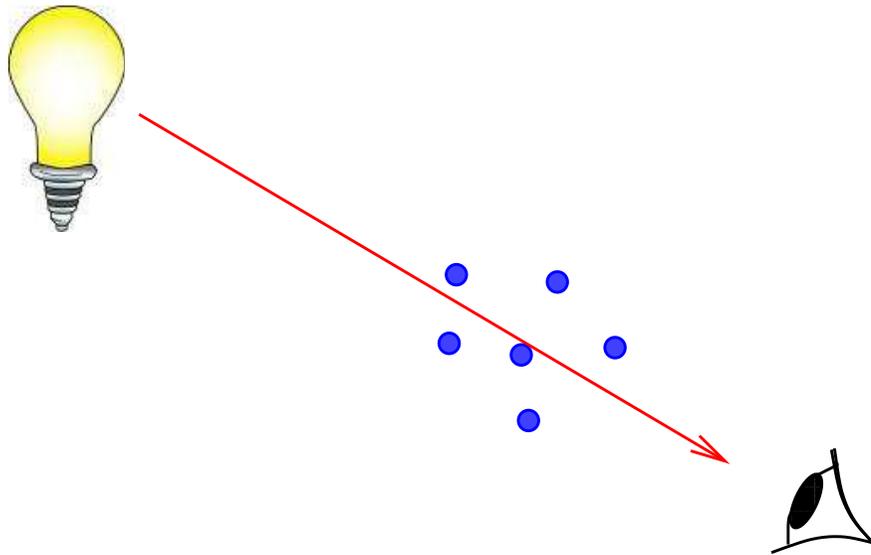
in collaboration with:

C. Hernández-Monteagudo & R. Sunyaev



Basic Idea

Metallicity evolution of the universe must leave its imprint on the Cosmic Microwave Background Radiation \Rightarrow through resonance scattering of the background photons by atoms, ions and molecules

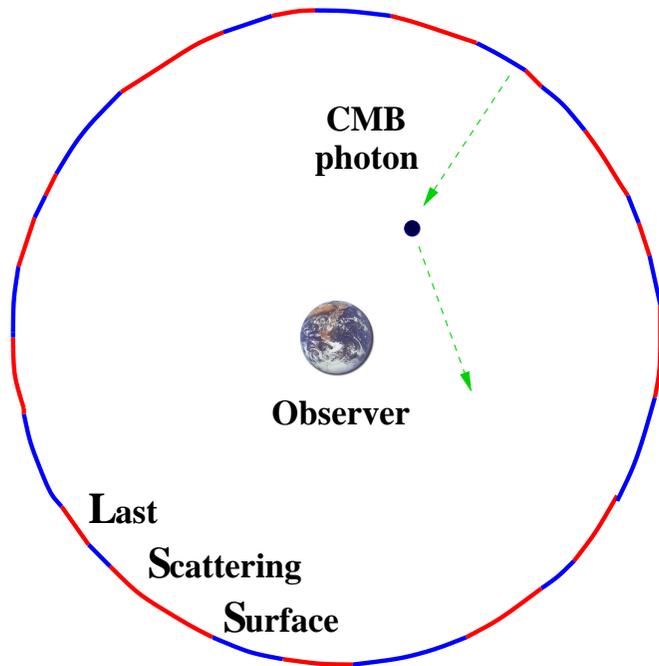


- Maoli, Melchiorri et al. 1994, 1996, ...
- Dubrovich 1977, 1993, ...
- de Bernardis et al. 1993
- Loeb & Zaldarriaga 2002
- Basu et al. 2004

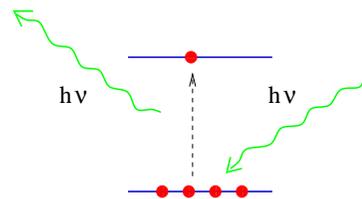


Suppression & generation

Resonant scattering partially erases the original temperature anisotropies of the CMB, but also generates new fluctuations at the epoch of scattering



$$\frac{\Delta T}{T_0}(\theta) = e^{-\tau_\nu} \left. \frac{\Delta T}{T_0}(\theta) \right|_{orig.} + \left. \frac{\Delta T}{T_0}(\theta) \right|_{new}$$

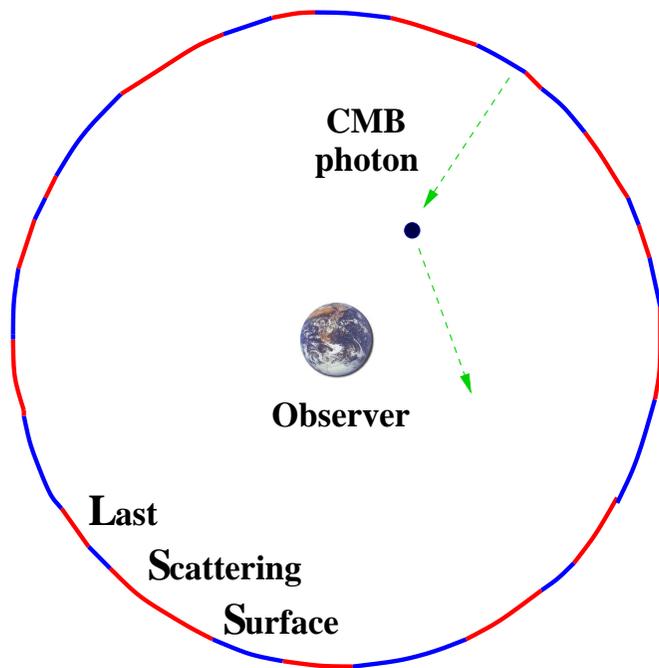


Scattering is *coherent*
in scatterer's frame



Suppression & generation

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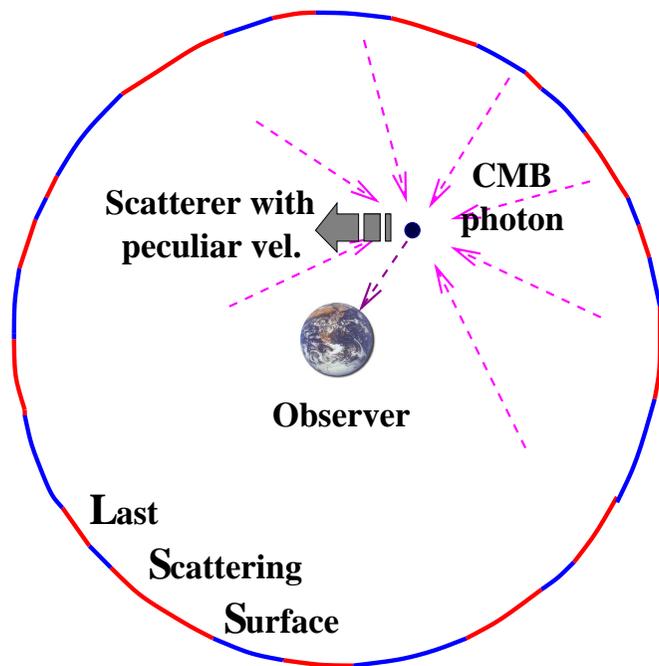
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“suppression (blurring)” & “generation”



Suppression & generation

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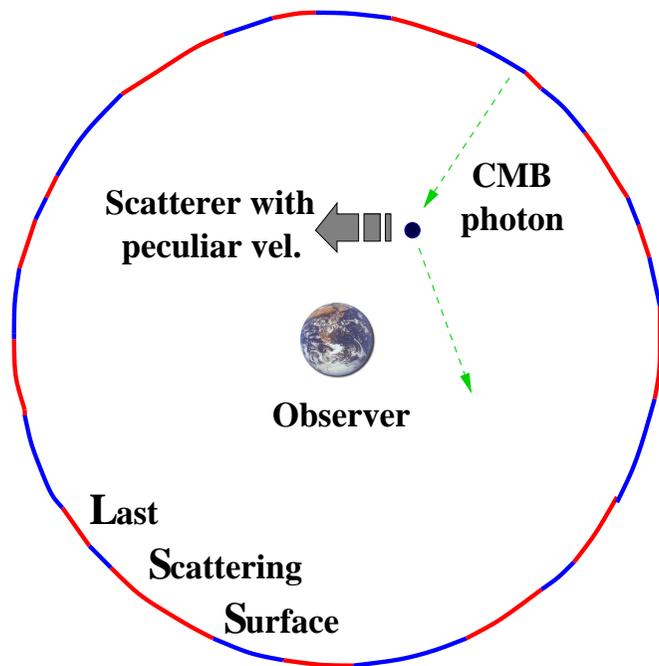
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Suppression & generation

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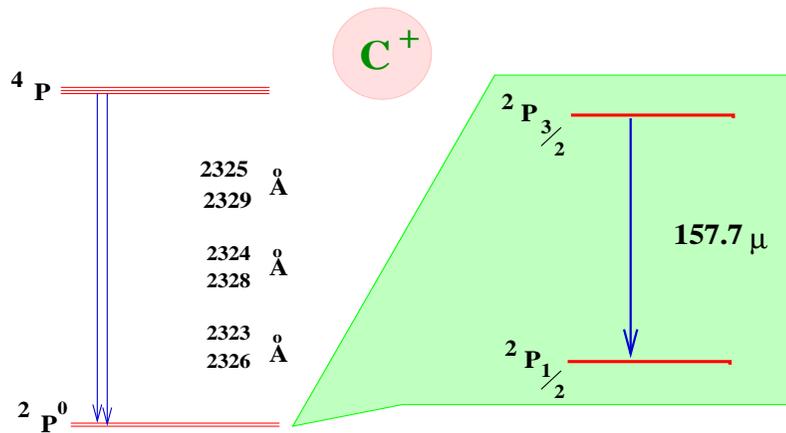
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$$\delta C_l = \tau_\nu \cdot \mathcal{C}_1 + \tau_\nu^2 \cdot \mathcal{C}_2 + \mathcal{O}(\tau_\nu^3)$$

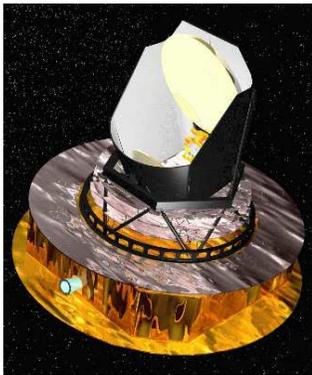


Fine-structure lines



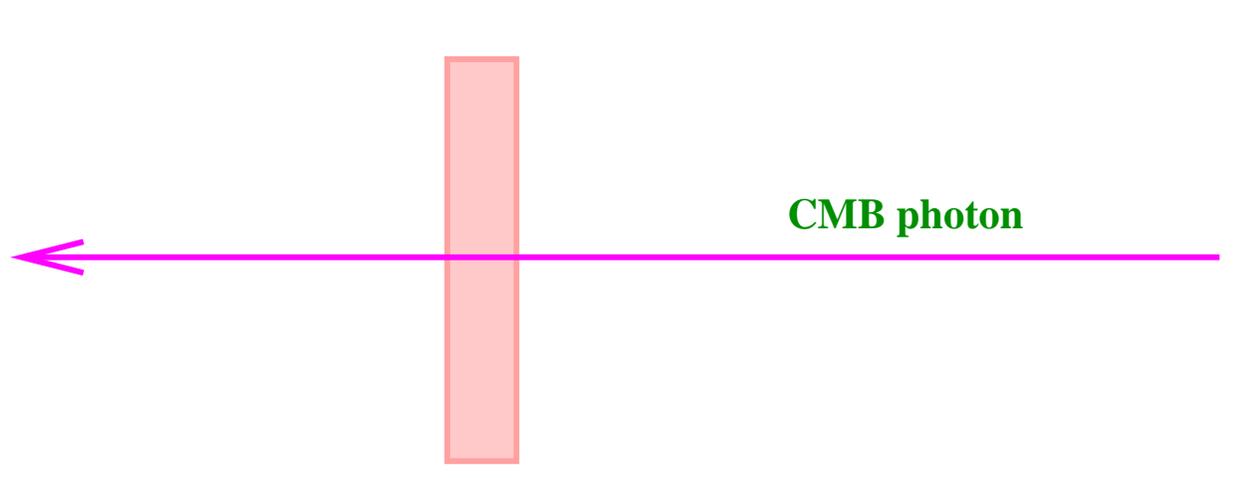
A-coefficient $\sim 10^{-6} \text{ s}^{-1}$
 \Rightarrow optical depth $\sim 10^{-5} - 10^{-7}$

143 GHz



Planck HFI

Bandwidth $\sim 25\%$



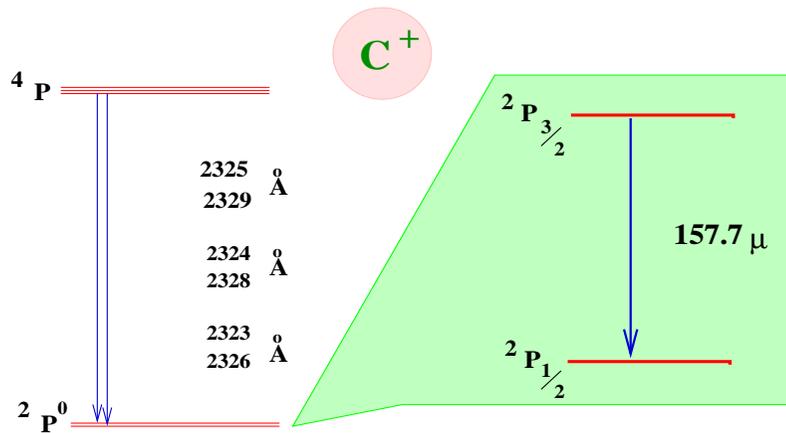
$z \sim 12$

LSS

at $z=1100$

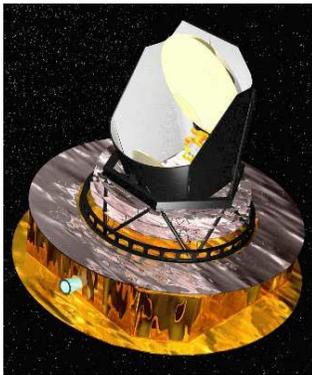


Fine-structure lines



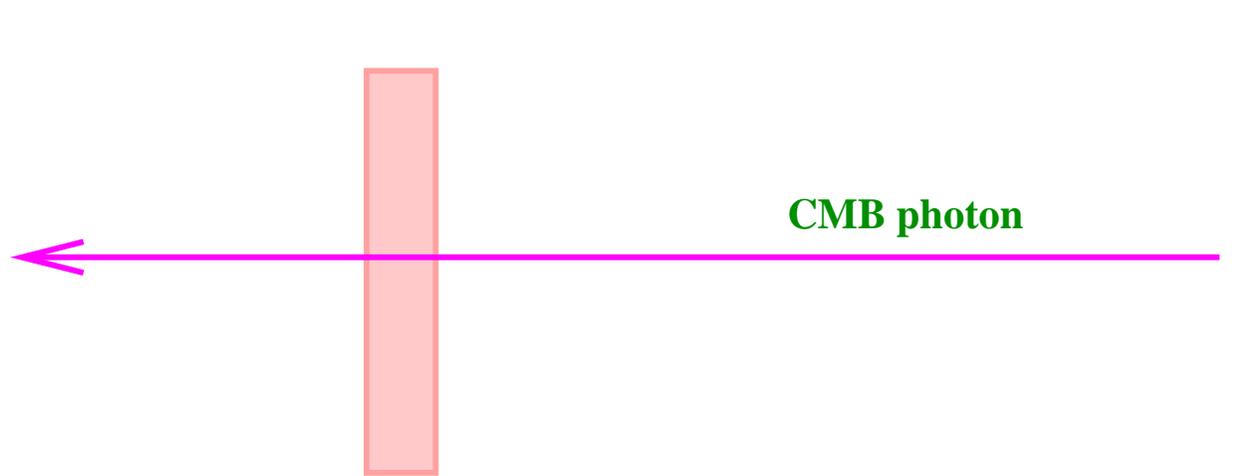
A-coefficient $\sim 10^{-6} \text{ s}^{-1}$
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217 GHz



Planck HFI

Bandwidth $\sim 25\%$



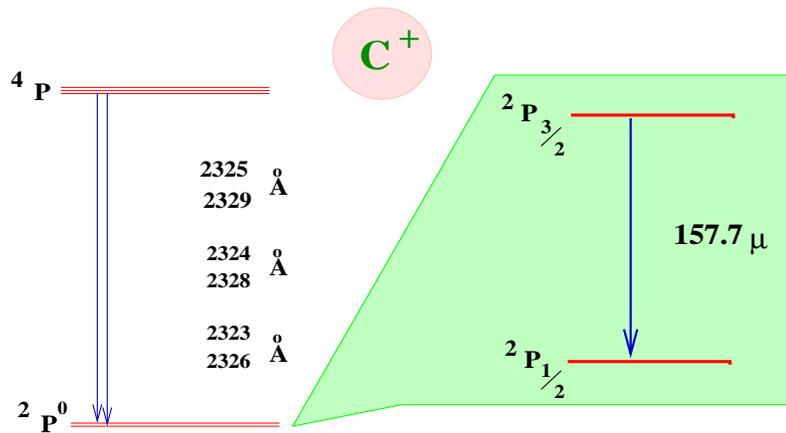
$z \sim 8$

LSS

at $z=1100$

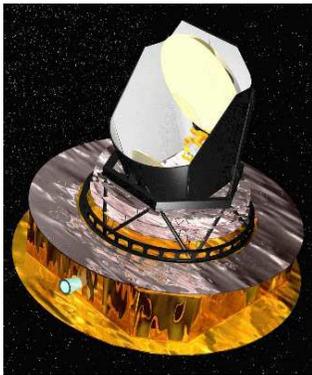


Fine-structure lines



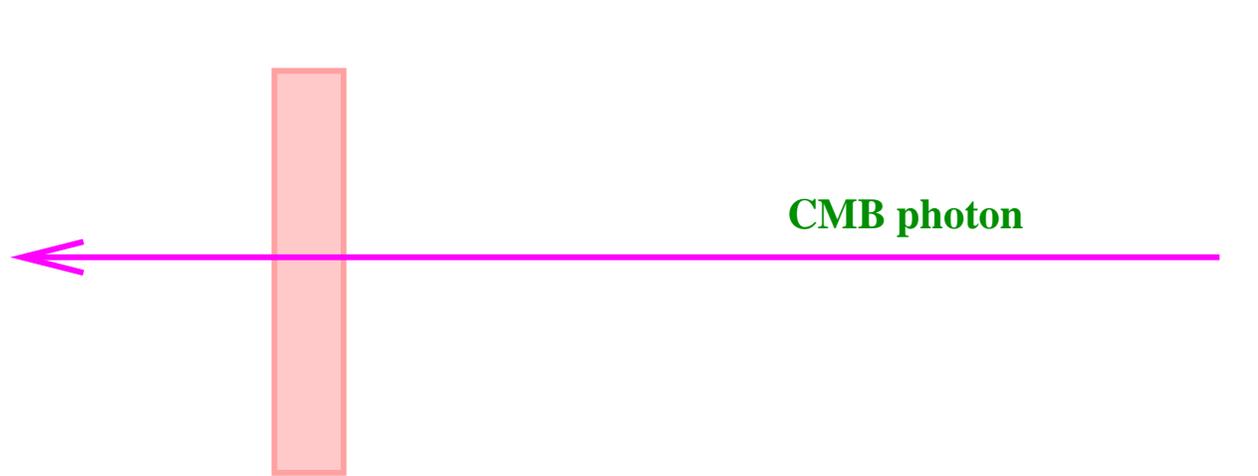
A-coefficient $\sim 10^{-6} \text{ s}^{-1}$
 \Rightarrow optical depth $\sim 10^{-5} - 10^{-7}$

353 GHz



Planck HFI

Bandwidth $\sim 25\%$



$z \sim 4$

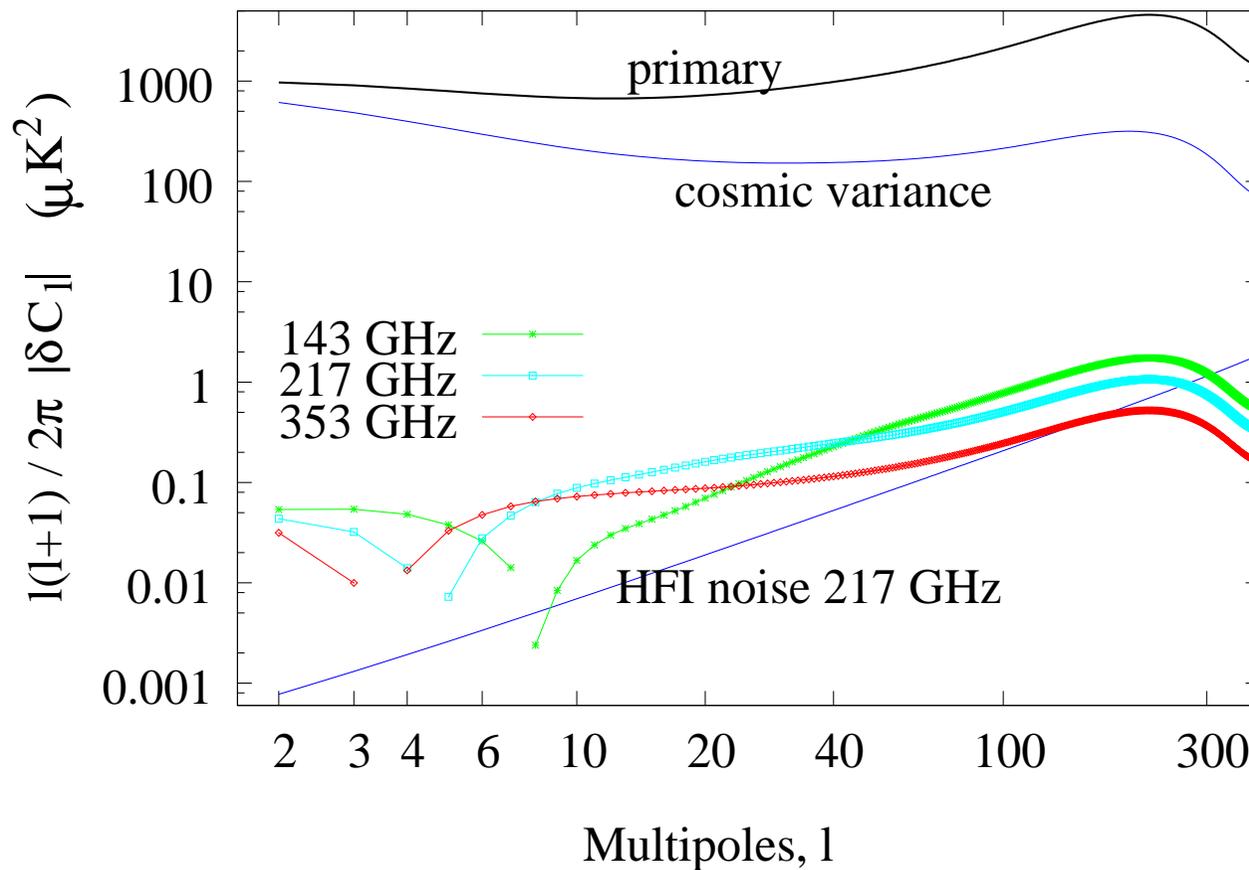
LSS

at $z=1100$

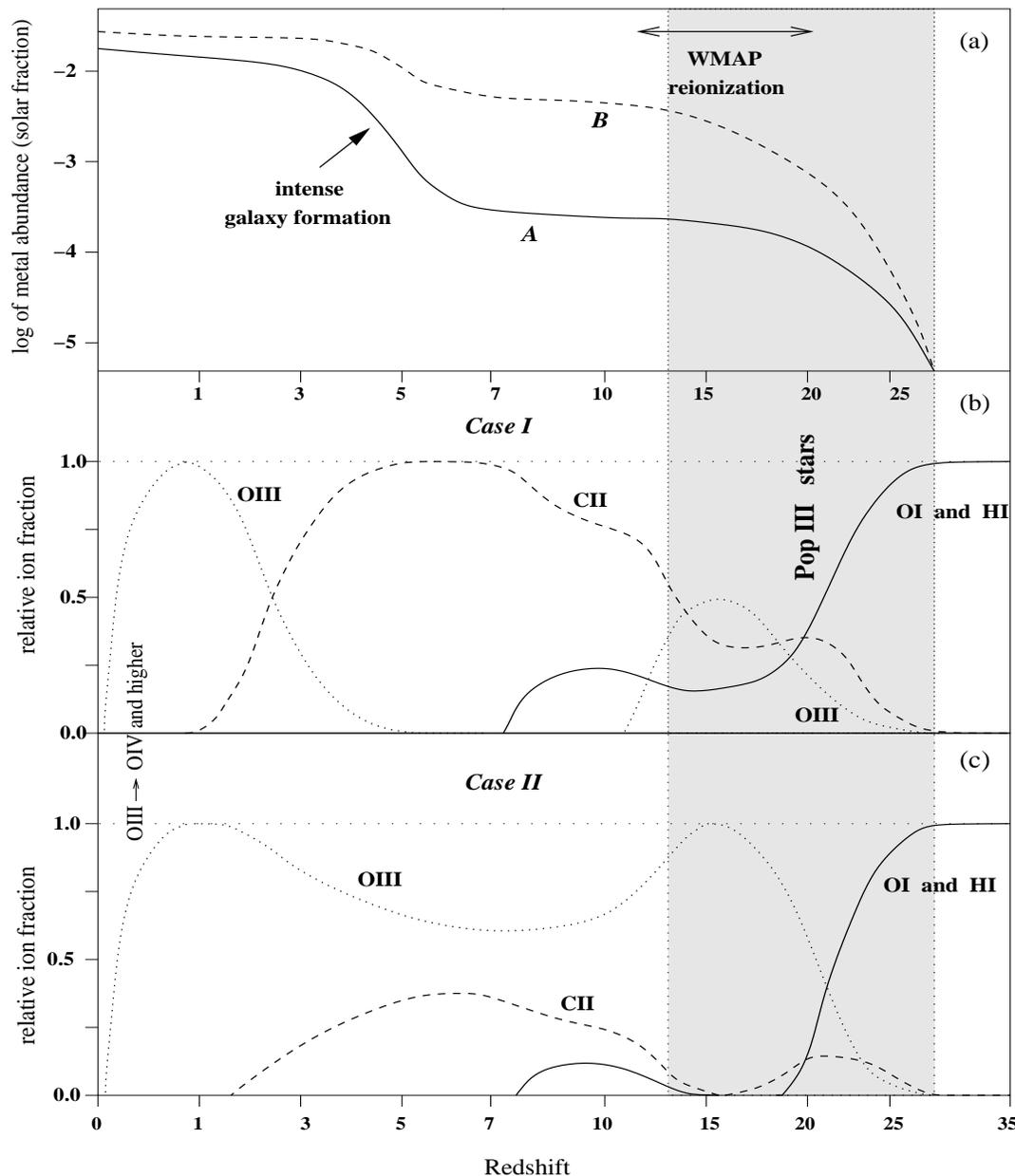


Frequency dependent δC_l

At small angular scales ($l \gtrsim 200$), we have simply $\delta C_l \simeq -2 \tau_{X_i} C_l^{prim.}$
making prediction of effect very easy for ground-based experiments!



Constraining enrichment histories



Different ionization & enrichment histories will produce distinct power spectrum distortions

For details see:

K. Basu,

C. Hernández-Monteagudo

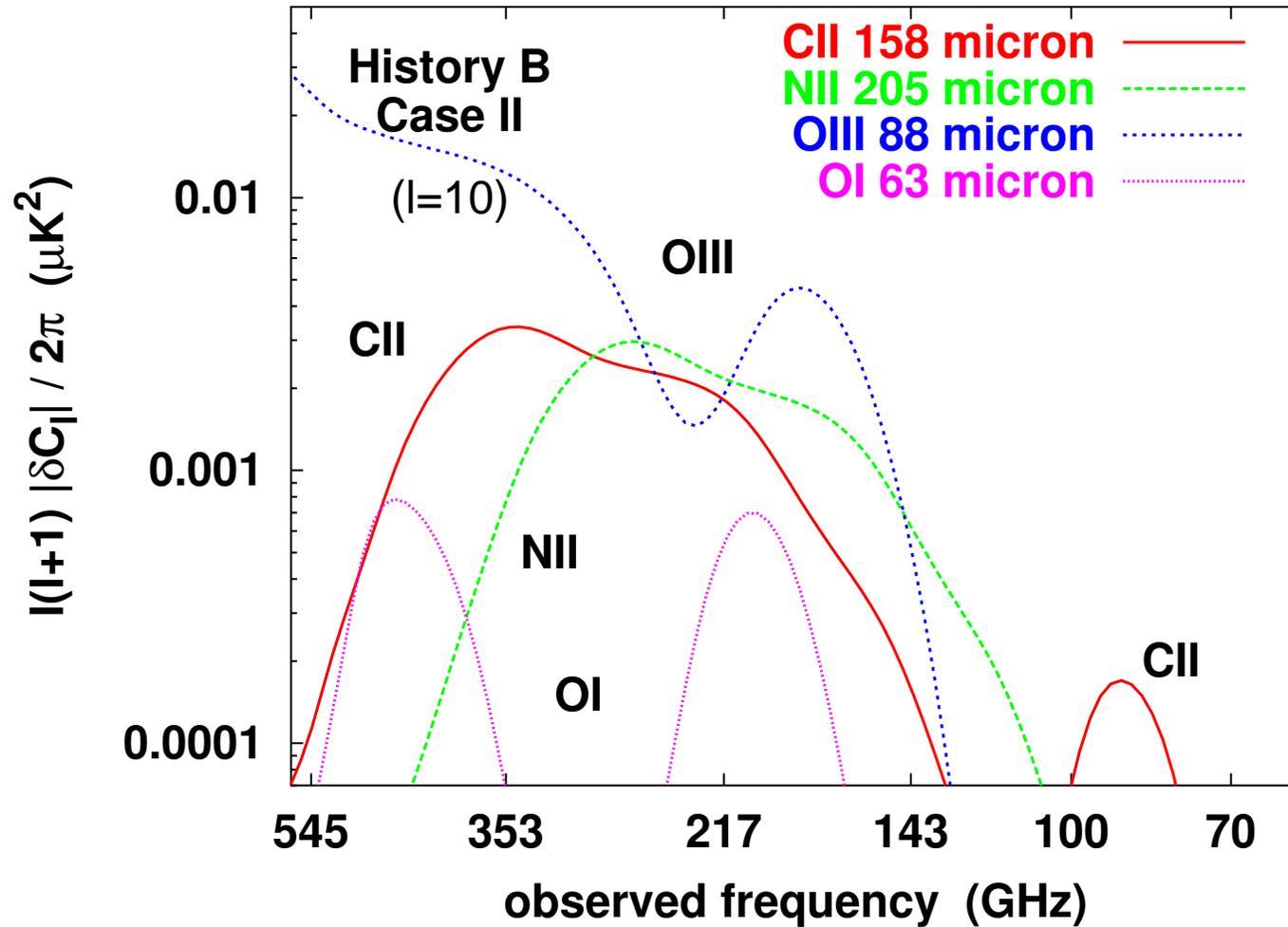
& R. Sunyaev

A&A, 2004, 416, 447

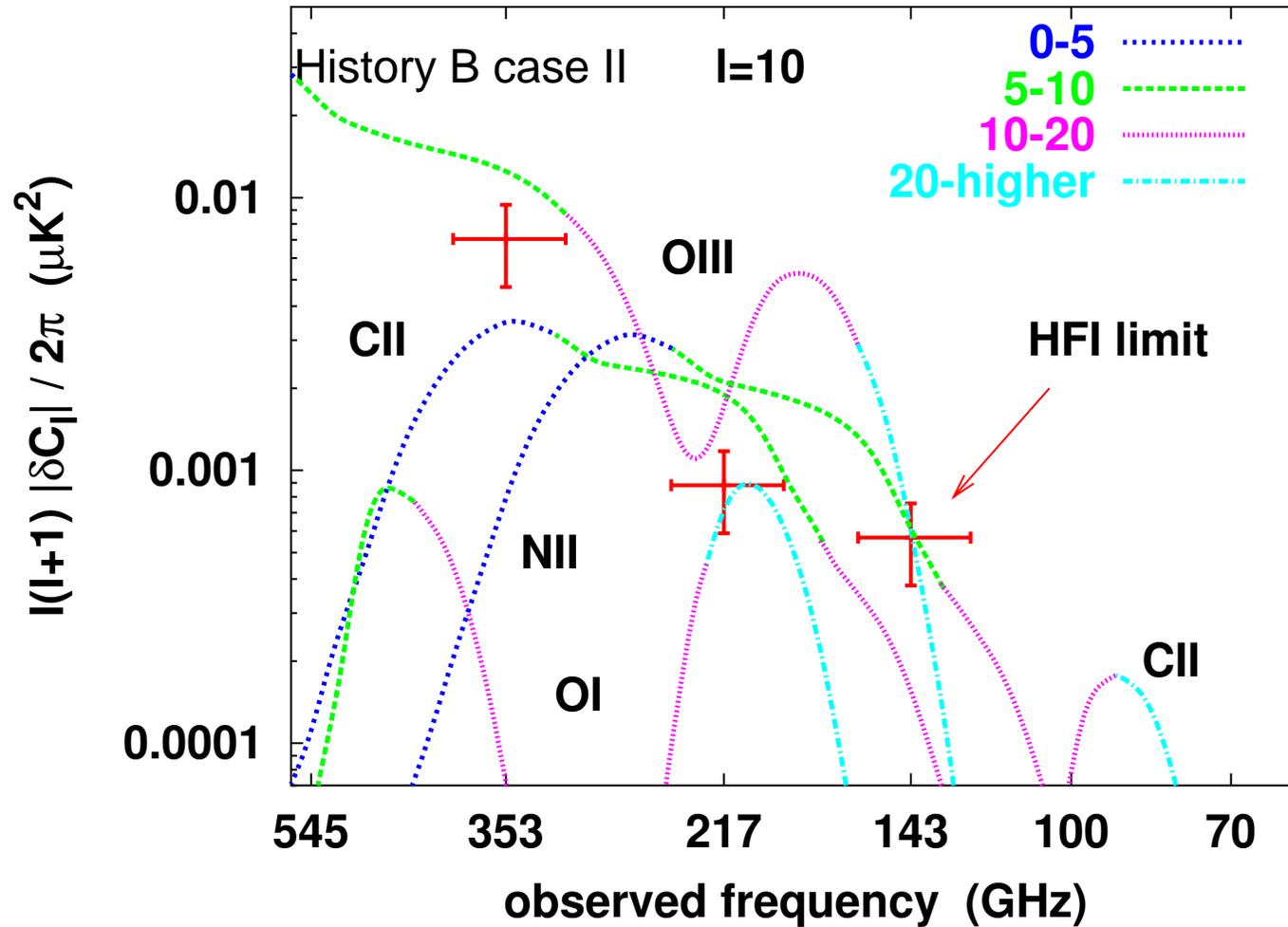
(detailed formalism, minimum abundances, foreground issues, ...)



Constraining enrichment histories



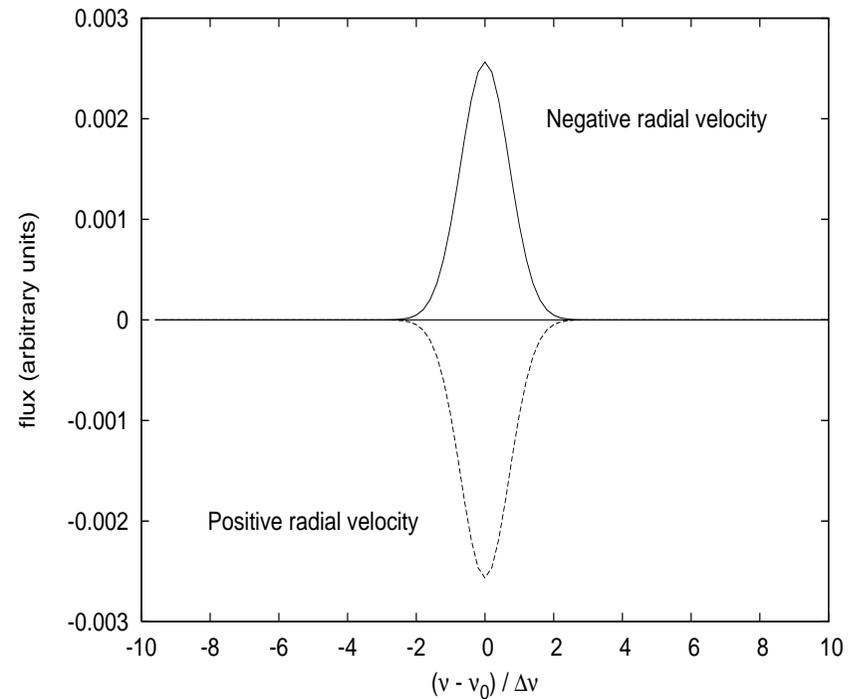
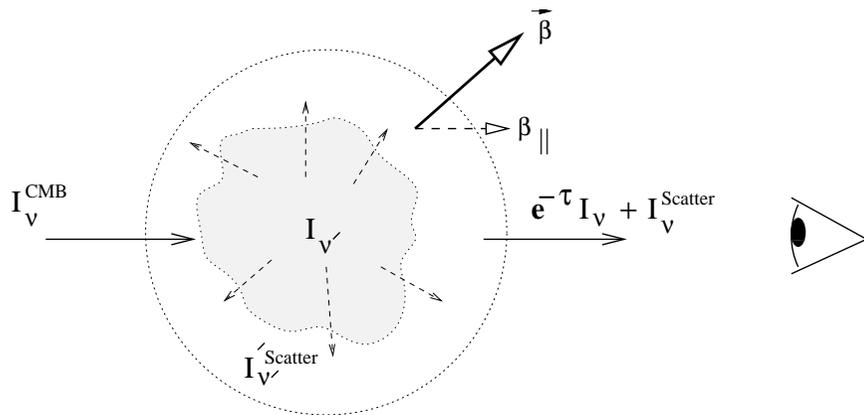
Constraining enrichment histories



Signal from resonant scattering

Very simple!

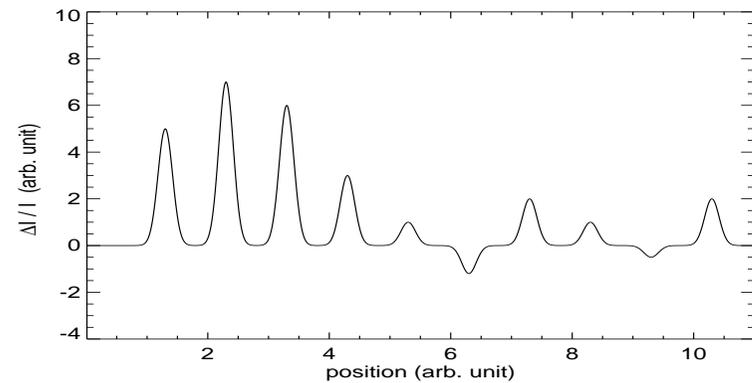
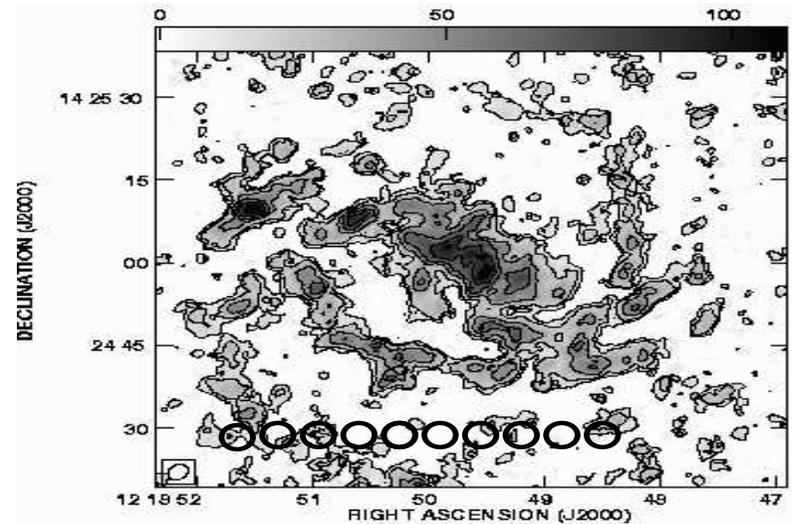
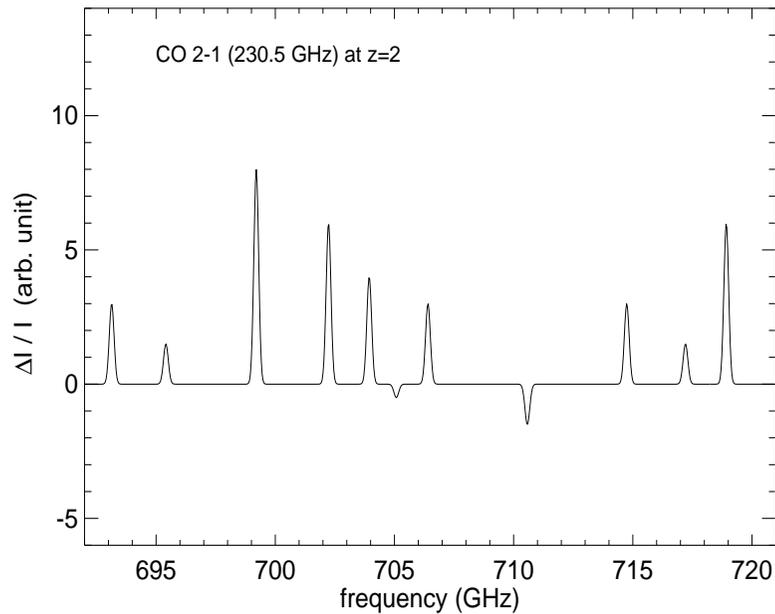
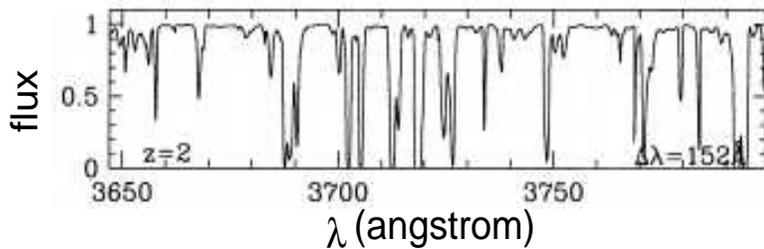
$$\frac{\Delta T}{T} = \tau_\nu \frac{v_{\text{rad}}}{c} \Rightarrow \delta I_\nu = I_0 \frac{x e^x}{e^x - 1} \frac{\Delta T}{T}; \quad x \equiv \frac{h\nu}{kT_{\text{CMB}}}$$



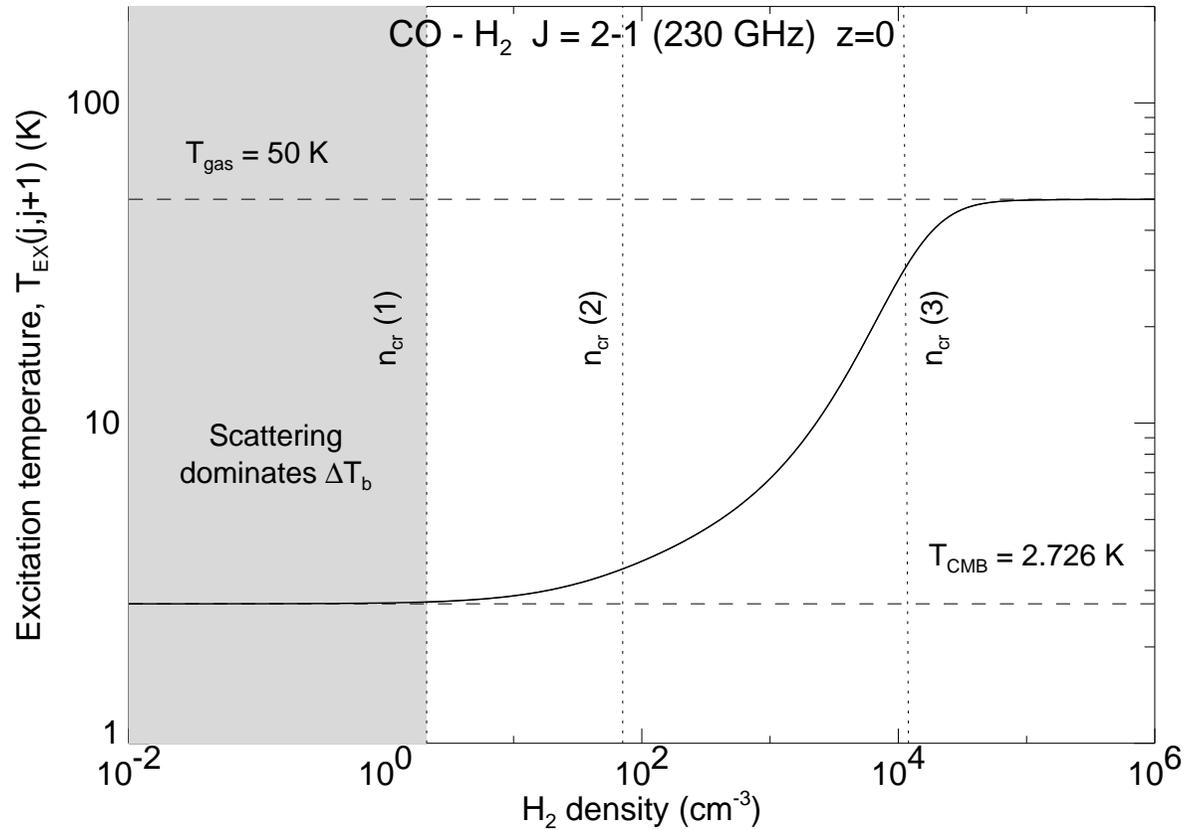
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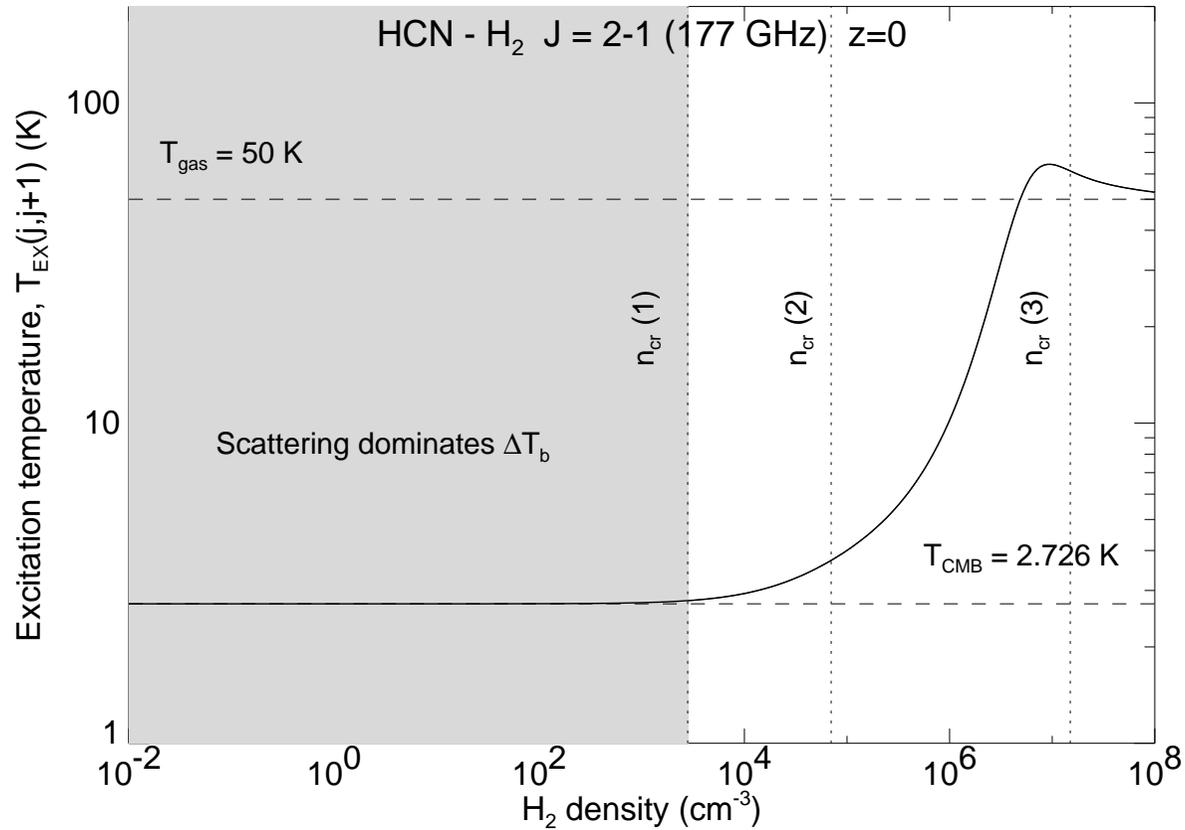
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Scattering vs. Emission

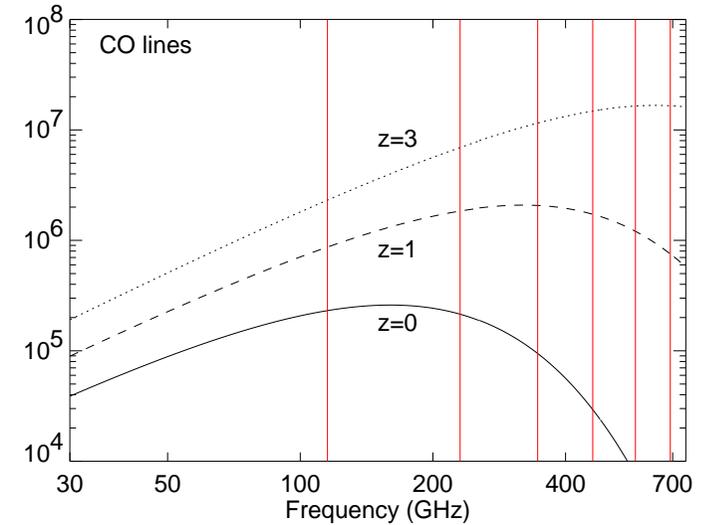
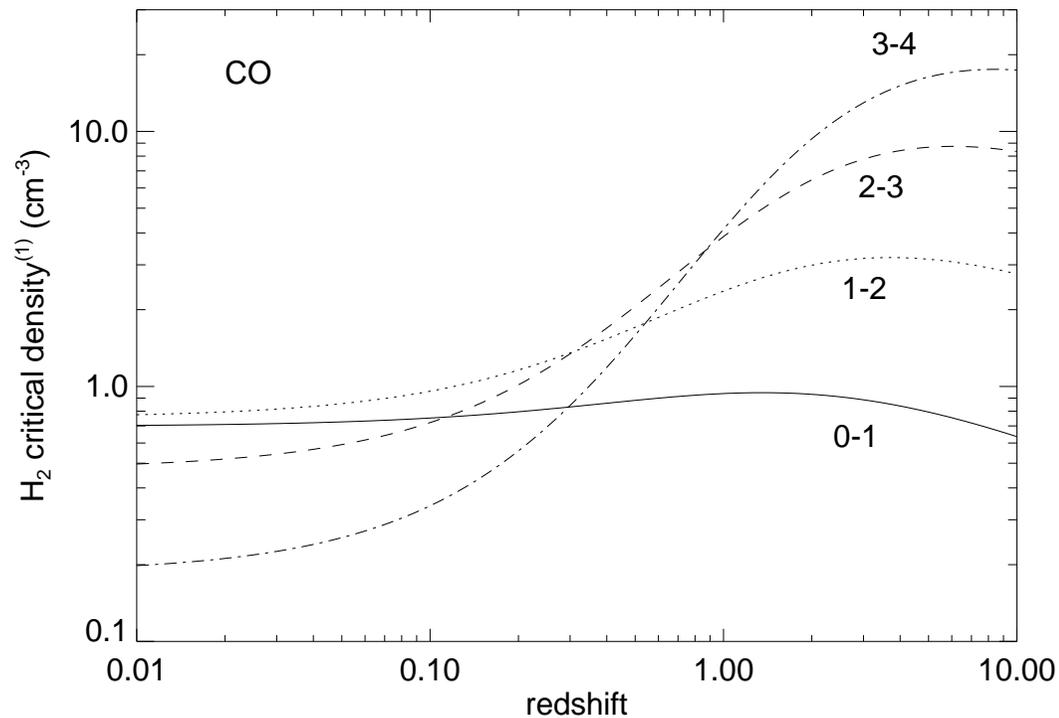


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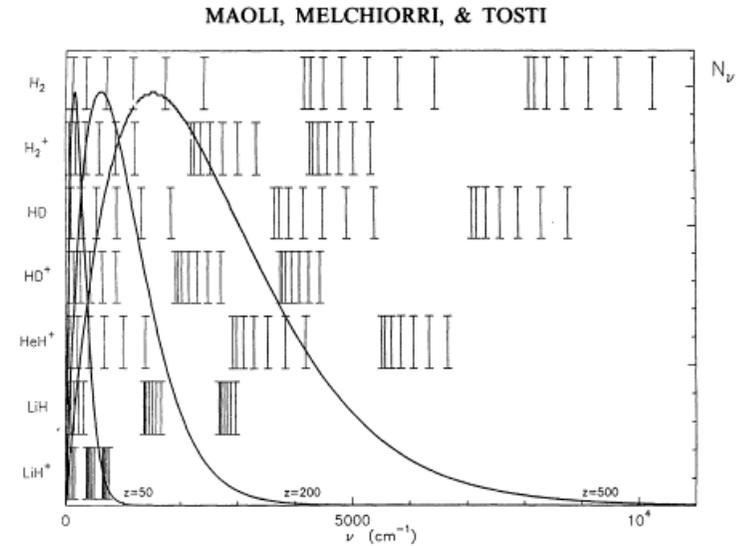
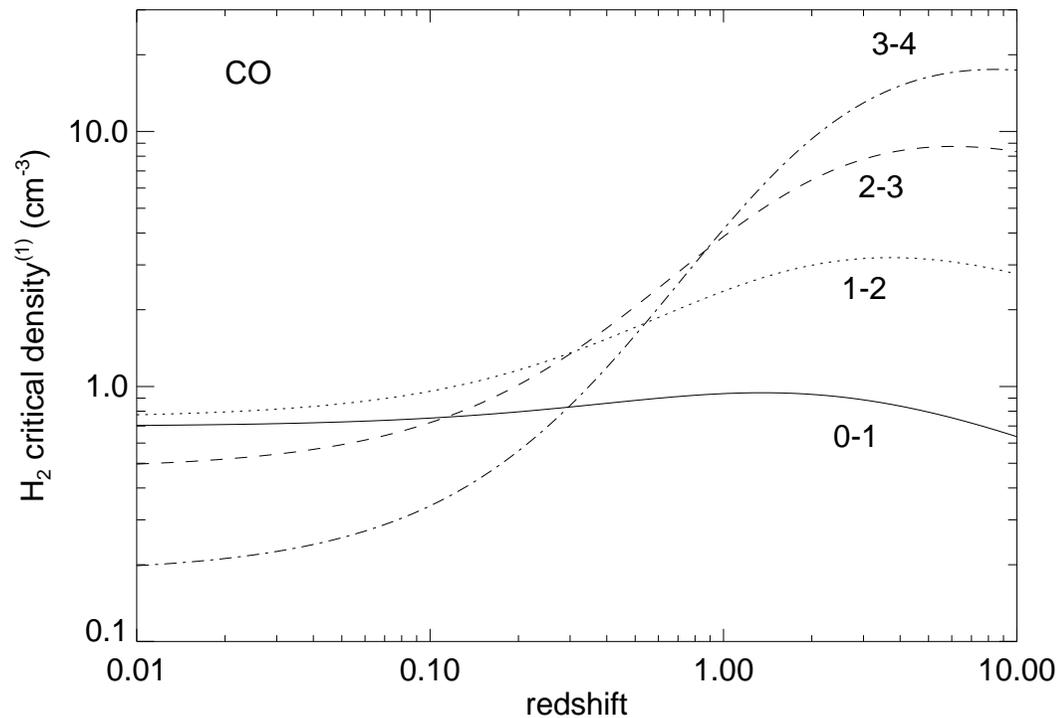
Scattering vs. Emission

At high- z CMB becomes more effective
higher rotational transitions show promise

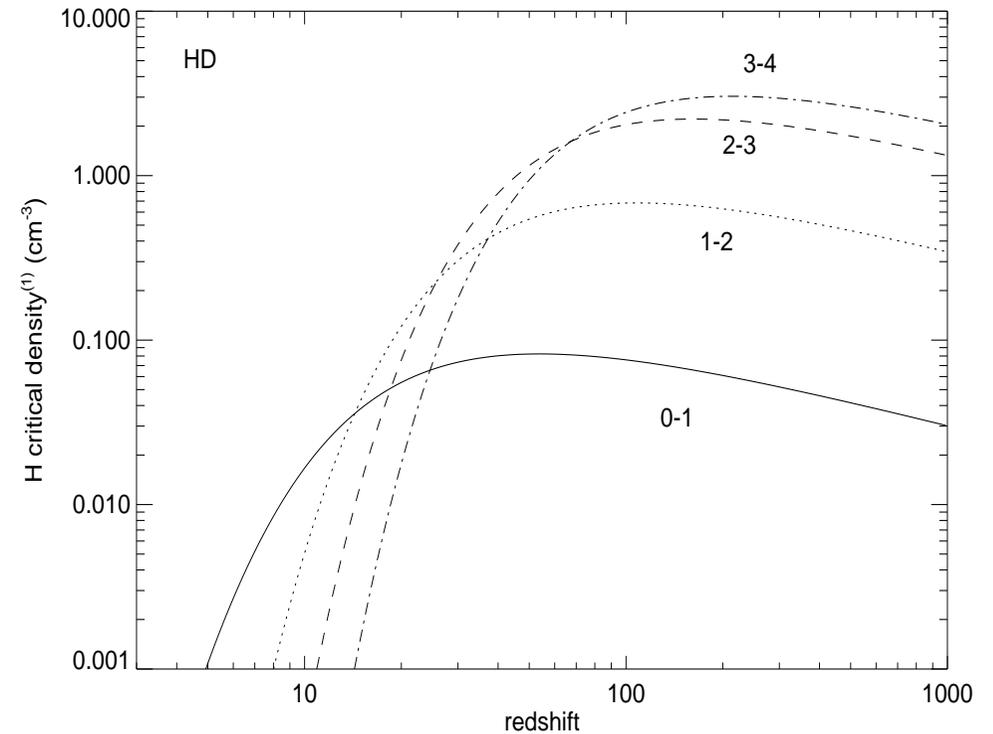
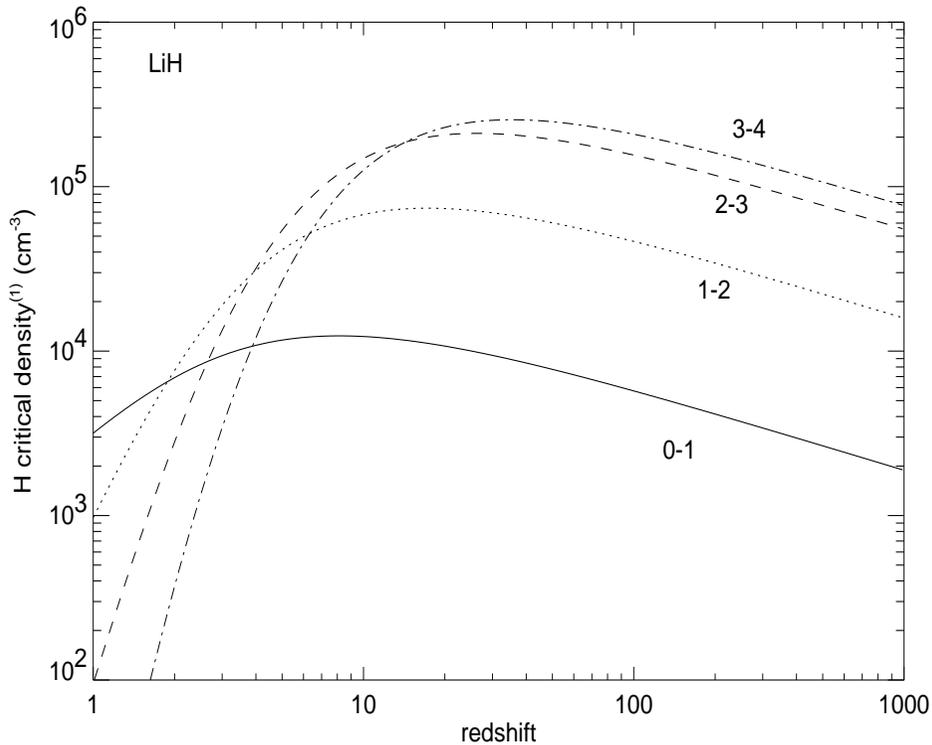


Scattering vs. Emission

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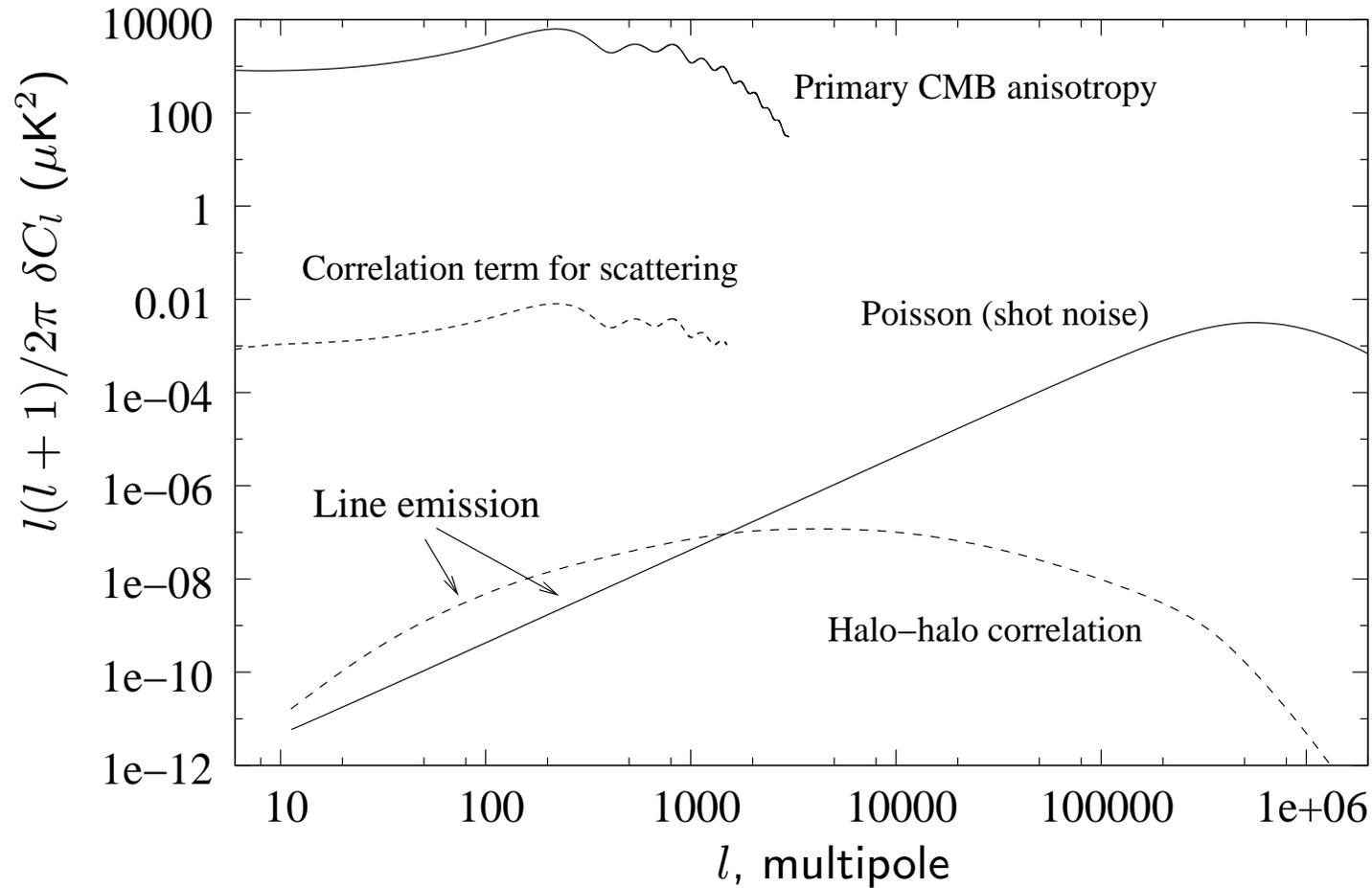
Primordial molecules



Extremely low abundance of LiH ($\sim 10^{-19}$ at $z \lesssim 100$) is a serious drawback.
Weak dipole moment of the HD molecule is another problem.



Primordial molecules



In a nutshell...

- The effect discussed by Dubrovich, Maoli, Melchiorri and others is very difficult to observe directly in individual objects
- Most promising place to look for the signature of resonant scattering is in the large scale coherent distortion of the CMB *power spectrum*, because effect is linear to the abundance and have very distinct spectral dependence
- Scattering caused by the fine-structure transitions of atoms and ions, produced by the first stars in the universe, can help to put tight constraints on the ionization and enrichment history with upcoming experiments like the Planck HFI

